

Part A

$$Q1. (i) \bar{E} = \frac{\sum_{n=0}^{\infty} E_n e^{-\frac{E_n}{KT}}}{\sum_{n=0}^{\infty} e^{-\frac{E_n}{KT}}}, \quad E_n = n\hbar\omega, \quad n=0,1,2,\dots$$

$$\text{Let } x = e^{-\frac{\hbar\omega}{KT}} \Rightarrow \bar{E} = \frac{\hbar\omega \sum_{n=0}^{\infty} nx^n}{\sum_{n=0}^{\infty} x^n}$$

$$\text{But } \sum_{n=0}^{\infty} x^n = \frac{1}{1-x} \quad n < 0 \text{ (i.e. } n \text{ is negative)}$$

if and only if ~~if and only if~~

Then

$$\frac{d}{dx} \sum_{n=0}^{\infty} x^n = \frac{d}{dx} \left( \frac{1}{1-x} \right) = \sum_{n=0}^{\infty} \frac{d}{dx} x^n$$

$$\Rightarrow \frac{d}{dx} (1-x)^{-1} = x^{-1} \sum_{n=0}^{\infty} nx^n$$

$$\therefore -\frac{1}{(1-x)^2} (-1) = x^{-1} \sum_{n=0}^{\infty} nx^n$$

$$\Rightarrow \sum_{n=0}^{\infty} nx^n = \frac{x}{(1-x)^2}$$

$$\text{Hence } \bar{E} = \hbar\omega \frac{\frac{x}{(1-x)^2}}{\frac{1}{(1-x)}} = \hbar\omega \frac{x}{1-x} = \hbar\omega \frac{e^{-\frac{\hbar\omega}{KT}}}{1 - e^{-\frac{\hbar\omega}{KT}}}$$

$$\therefore \bar{E} = \frac{\hbar\omega}{e^{\frac{\hbar\omega}{KT}} - 1} \quad (\text{per atomic oscillator, per degree of freedom}).$$

(ii)

$\Rightarrow$  For 3-degrees of freedom,

$$\bar{E} = \frac{3\hbar\omega}{e^{\frac{\hbar\omega}{KT}} - 1}$$

Therefore, per mole

2.

$$E = N_A \bar{E} / 3 \text{ degrees} = \frac{3 N_A \hbar \omega}{e^{\frac{\hbar \omega}{kT}} - 1}$$

Then

$$C_V = \left. \frac{\partial E}{\partial T} \right|_V = 3 N_A \cdot \hbar \omega \left[ \frac{\partial}{\partial \left( e^{\frac{\hbar \omega}{kT}} - 1 \right)} \left( e^{\frac{\hbar \omega}{kT}} - 1 \right)^{-1} \cdot \frac{\partial}{\partial T} \left( e^{\frac{\hbar \omega}{kT}} - 1 \right) \right]$$

$$= 3 N_A \cdot \hbar \omega \left[ - \left( e^{\frac{\hbar \omega}{kT}} - 1 \right)^{-2} \cdot \frac{\partial}{\partial \left( \frac{\hbar \omega}{kT} \right)} \left( e^{\frac{\hbar \omega}{kT}} - 1 \right) \cdot \frac{\partial}{\partial T} \left[ \left( \frac{\hbar \omega}{k} \right) T^{-1} \right] \right]$$

$$= - 3 N_A \cdot \hbar \omega \cdot \frac{e^{\frac{\hbar \omega}{kT}}}{\left( e^{\frac{\hbar \omega}{kT}} - 1 \right)^2} \cdot \left( \frac{\hbar \omega}{k} \right) \left( - \frac{1}{T^2} \right)$$

$$= 3 k N_A \cdot \frac{\left( \frac{\hbar \omega}{kT} \right)^2}{\left( e^{\frac{\hbar \omega}{kT}} - 1 \right)^2} \cdot \frac{e^{\frac{\hbar \omega}{kT}}}{\left( e^{\frac{\hbar \omega}{kT}} - 1 \right)^2}$$

But  $R = k N_A$  and  $\omega \rightarrow \omega_E$ , the Einstein frequency

$$\Rightarrow C_V = 3 R \left( \frac{\hbar \omega_E}{kT} \right)^2 \cdot \frac{e^{\frac{\hbar \omega_E}{kT}}}{\left( e^{\frac{\hbar \omega_E}{kT}} - 1 \right)^2}$$

which is the required expression.

Q2. At the low-temperature limit of  $C_V$  in the Debye model,

3.

$$C_V \approx \frac{12}{5} \pi^4 k \left( \frac{T}{\theta_D} \right)^3 \Rightarrow C_V = \alpha \left( \frac{T}{\theta_D} \right)^3$$

$$\therefore \alpha = \frac{C_V}{\left( \frac{T}{\theta_D} \right)^3} = C_V \left( \frac{\theta_D}{T} \right)^3$$

Using KCl at 5K,  $\alpha = \left( \frac{230}{5} \right)^3 (3.8 \times 10^{-2}) = 3.6988 \times 10^3 \text{ JK}^{-1} \text{ mol}^{-1}$

$$\Rightarrow C_V = 3.6988 \times 10^3 \left( \frac{T}{\theta_D} \right)^3$$

Hence for NaCl at 5K,  $C_V = 3.6988 \times 10^3 \left( \frac{5}{310} \right)^3$   
 $\approx \underline{\underline{1.552 \times 10^{-2} \text{ JK}^{-1} \text{ mol}^{-1}}}$

Similarly, for KCl at 2K,  $C_V = 3.6988 \times 10^3 \left( \frac{2}{230} \right)^3$   
 $\approx \underline{\underline{2.432 \times 10^{-3} \text{ JK}^{-1} \text{ mol}^{-1}}}$

Q3. Let the mass of each particle be  $m$ . Then

$$m \ddot{u}_r = F_r \quad \text{--- (1)}$$

But  $u_x = A e^{i(kx - \omega t)} \Rightarrow u_r = A e^{i(kra - \omega t)}$ , where  $\alpha = ra$ .

$$\Rightarrow \frac{d^2 u_r}{dt^2} = -\omega^2 u_r \quad \text{or} \quad \ddot{u}_r = -\omega^2 u_r \quad \text{--- (2)}$$

(2) into (1) gives

$$-m\omega^2 u_r = F_r \quad \text{--- (3)}$$

From Hooke's law, the force on the  $r$ -th atom is

$$F_r = \mu(u_{r+1} - u_r) - \mu(u_r - u_{r-1})$$

where  $\mu$  is the atomic force constant (i.e. elasticity at inter-atomic level)

$$\therefore F_r = \mu [u_{r+1} + u_{r-1} - 2u_r] \quad \text{--- (4)}$$

But (3) = (4)

$$\Rightarrow \omega^2 = + \frac{\mu}{m u_r} \left[ 2 - \frac{u_{r+1} + u_{r-1}}{u_r} \right] \quad \text{--- (5)}$$

$$\text{But } u_r = A e^{i(kra - \omega t)} = A e^{iakra} e^{-i\omega t}$$

$$u_{r+1} = A e^{i[k(r+1)a - \omega t]} = A e^{-i\omega t} e^{ikr} e^{ika}$$

$$u_{r-1} = A e^{i[k(r-1)a - \omega t]} = A e^{-i\omega t} e^{ikr} e^{-ika}$$

Hence (5) becomes

$$\omega^2 = \frac{\mu}{m} \left( 2 - e^{ika} - e^{-ika} \right)$$

$$\text{But } e^{\pm i\phi} = \cos\phi \pm i \sin\phi \Rightarrow \cos\phi = \frac{e^{i\phi} + e^{-i\phi}}{2}$$

where if  $\phi = ka$  then

$$\omega^2 = \frac{\mu}{m} (2 - 2\cos ka)$$

$$= \frac{2\mu}{m} (1 - \cos ka)$$

Furthermore,  $\cos 2\phi = 1 - 2\sin^2\phi$

$$\therefore \frac{1 - \cos 2\phi}{2} = \sin^2\phi$$

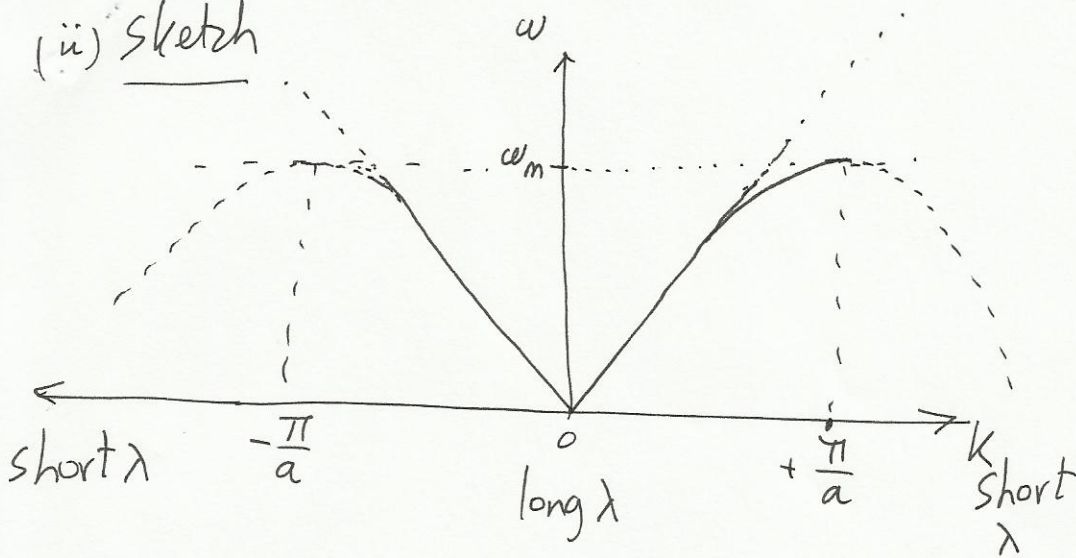
$$\text{Let } 2\phi = ka \Rightarrow \phi = \frac{ka}{2}$$

$$\therefore \omega^2 = \frac{4\mu}{m} \left( \frac{1 - \cos ka}{2} \right) = \frac{4\mu}{m} \sin^2\left(\frac{ka}{2}\right)$$

$$\Rightarrow \omega = \pm 2 \sqrt{\frac{\mu}{m}} \sin\left(\frac{ka}{2}\right) \text{ or } \omega = \pm \omega_m \sin\left(\frac{ka}{2}\right)$$

where  $\omega_m = 2\sqrt{\frac{\mu}{m}}$ .

(ii) Sketch



$$k = \frac{2\pi}{\lambda}$$

( $\lambda \rightarrow 0 = \text{short } \lambda$   
 $\Rightarrow k \text{ is large}$ )

$$1 = \sin \frac{ka}{2} \equiv \text{maximum, when}$$

$$\frac{ka}{2} = \pm \frac{\pi}{2} \therefore k = \pm \frac{\pi}{a}$$

$$0 = \sin \frac{ka}{2}, \text{ when } \frac{ka}{2} = \pm \pi \Rightarrow k = \pm \frac{2\pi}{a}$$

Q4.  $\omega_D = v(6\pi^2 n)^{\frac{1}{3}} = v(6\pi^2)^{\frac{1}{3}} n^{\frac{1}{3}}$

$$\omega_D \hbar = k T_D \Rightarrow \frac{k T_D}{\hbar} = v(6\pi^2)^{\frac{1}{3}} \left( \frac{M_A}{M} \right)^{\frac{1}{3}}$$

$$\Rightarrow T_D = \frac{\hbar}{k} v(6\pi^2 M_A)^{\frac{1}{3}} \left( \frac{\rho}{M} \right)^{\frac{1}{3}} = \beta \downarrow v \left( \frac{\rho}{M} \right)^{\frac{1}{3}} \text{ const.}$$

$$\Rightarrow \frac{T_D^{(Au)}}{T_D^{(Cu)}} = \frac{v^{(Au)} \left[ \frac{\rho^{(Au)}}{M^{(Au)}} \right]^{\frac{1}{3}}}{v^{(Cu)} \left[ \frac{M^{(Cu)}}{\rho^{(Cu)}} \right]^{\frac{1}{3}}}$$

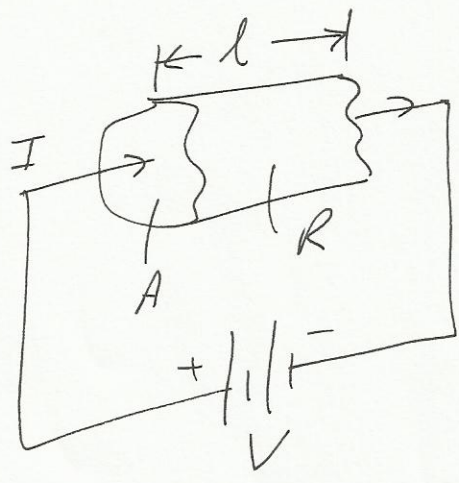
$$\Rightarrow T_D^{(Au)} = \frac{v^{(Au)}}{v^{(Cu)}} \left[ \frac{\rho^{(Au)}}{\rho^{(Cu)}} \cdot \frac{M^{(Cu)}}{M^{(Au)}} \right]^{\frac{1}{3}} T_D^{(Cu)}$$

$$= \left( \frac{2.1}{3.8} \right) \left[ \frac{19.3}{8.9} * \frac{63.5}{197} \right]^{\frac{1}{3}} 348 \text{ Kelvin}$$

$$\therefore \underline{\underline{T_D^{(Au)} = 170.7 \text{ K}}}$$

Part B.

5. (i) Ohm's Law  $\Rightarrow V = I \cdot R$ .



$$R = \rho \frac{L}{A}$$

(electric field)  $E = \frac{V}{L} \Rightarrow E \cdot L = V$

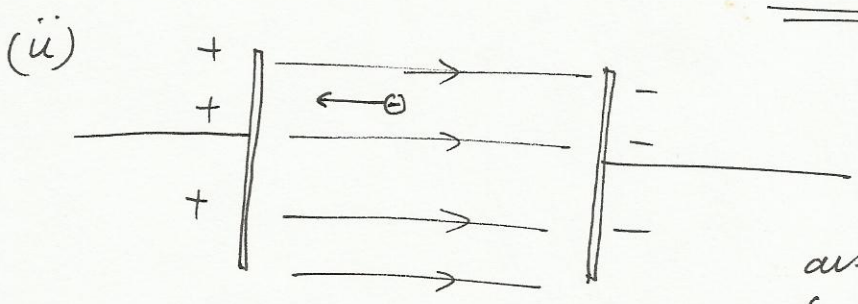
$$J = \frac{I}{A} \Rightarrow I = J \cdot A$$

$\therefore$  Since  $V = I \cdot R$

$$(E \cdot L) = (J \cdot A) \cdot \left(\rho \frac{L}{A}\right)$$

$$\Rightarrow E = \rho J \quad \text{or} \quad J = \frac{1}{\rho} E$$

$$\underline{\underline{J = \sigma E}}$$



$$F = qE.$$

drift force due to external potential  $= (-eE)$

average force due to electron collisions  $F_c = m \frac{v_d}{\tau}$

net force on electrons  $F_n$

$$F_n = F_d - F_c$$

$$\Rightarrow m \frac{dv}{dt} = -eE - m \frac{v_d}{\tau}$$

At equilibrium,  $F_n = 0$

7.

$$\Rightarrow v_d = -\frac{e\tau}{m} \mathcal{E}$$

But  $J = -nev_d$  { where  $n =$  number of electrons per unit volume.

To see this, consider a time interval of  $\Delta t$  sec. Hence volume swept out by cross-section area is  $A v_d \Delta t$ , and number of electrons therein is  $(n A v_d \Delta t)$ .

Hence charge delivered in 1 sec is

$(-n e A v_d)$  and the current is

$$I = \frac{\Delta \text{charge}}{\text{time}} = \frac{-n e A v_d \Delta t}{\Delta t} = -n e A v_d$$

By definition

$$J = \frac{I}{A} = -n e v_d$$

$$\Rightarrow J = -n e \left( -\frac{e\tau}{m} \right) \mathcal{E}$$

$$\therefore J = \frac{n e^2 \tau}{m} \mathcal{E} \quad \text{which also equals } (\sigma \mathcal{E})$$

$$\therefore \sigma = \frac{n e^2 \tau}{m}$$

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(iii)  $M = 63.5 \text{ g}$   
 $\rho_{\text{den.}} = 8.95 \text{ g cm}^{-3}$   
 $\rho_{\text{res}} = 1.55 \times 10^{-8} \text{ } (\Omega \text{ m})$   
 $Z_V = 1.$

(a)  $n = \frac{Z_V \rho_{\text{den.}} N_A}{M} = \frac{1 \times 8.95 \text{ g cm}^{-3} \times 6.022 \times 10^{23} \text{ mol}^{-1}}{63.5 \text{ g mol}^{-1}}$   
 $= \underline{8.488 \times 10^{22} \text{ electrons per cm}^3}$   
 or  $n = \underline{8.488 \times 10^{28} \text{ per m}^3}$  ( $1 \text{ cm}^3 = 10^{-6} \text{ m}^3$ )

(b)  $\sigma = \frac{ne^2 \tau}{m} = \frac{1}{\rho_{\text{res.}}} \therefore \tau = \frac{m}{\rho_{\text{res.}} ne^2}$   
 $= \frac{9.11 \times 10^{-31} \text{ kg.}}{1.55 \times 10^{-8} \Omega \text{ m.} \times 8.488 \times 10^{28} \text{ m}^{-3} (1.6022 \times 10^{-19} \text{ C})^2}$   
 $= 2.697 \times 10^{-14} \text{ s}$

(c)  $E_F = \frac{\hbar^2}{2m} (3\pi^2 n)^{\frac{2}{3}}$   
 $= \frac{(6.626 \times 10^{-34})^2}{(2\pi)^2 \times 9.11 \times 10^{-31}} \left( 3\pi^2 \times 8.488 \times 10^{28} \right)^{\frac{2}{3}} \text{ Joules}$   
 $= 1.1283 \times 10^{-18} \text{ J}$   
 $= \underline{7.04 \text{ eV}}$   
 { but  $1 \text{ eV} = 1.6022 \times 10^{-19} \text{ J}$   
 $\therefore 1 \text{ J} = \frac{1}{1.6022 \times 10^{-19}} \text{ eV}$

(see also Table 4.11 for  $E_F^{(\text{Cu})}$ ).

(d) Fermi-velocity,  $V_F$  is  $\sqrt{\frac{2E_F}{m}} = \sqrt{\frac{2 \times 1.1283 \times 10^{-18}}{9.11 \times 10^{-31}}} \text{ m/s} = 1.574 \times 10^6 \text{ m/s}$   
 $\therefore \text{Mean free path} = \ell_F = V_F \tau = 2.697 \times 10^{-14} \times 1.574 \times 10^6 \text{ m} = 4.245 \times 10^{-8} \text{ m} = \underline{42.45 \text{ nm.}}$

## 6. The Drude theory (a.k.a the Classical theory).

9.

### Successes

- (i) correctly derives Ohm's law
- (ii) correctly predicts the dependence of resistivity on temperature
- (iii) explains the direct link between thermal conductivity and electrical resistivity.

### Failures

- (i) fails to explain the conduction electron contribution to heat capacity (thermal conductivity)
- (ii) also fails to explain the contribution of conduction electrons to magnetic susceptibility
- (iii) cannot explain why the mean free paths are so much longer than the lattice parameter.

## The free electron gas theory

### Failures

- (i) Its greatest failure is considered to be the sign of the Hall coefficient which can be both positive ~~and~~ or negative. The theory expects that  $R_H$  will always be negative.
- (ii) It expects that the number of conduction electrons in unit volume will increase with valency. This is not always so.
- (iii) It always predicts a spherical Fermi-surface whereas measurements often indicate it can be non-spherical.

(iii) Na;  $E_F = 3.23\text{eV}$

(a) Fermi-energy - is the energy level of the highest or most energetic of the ground states of a material at zero Kelvin.

(b)  $E_F = kT_F \Rightarrow T_F = \frac{E_F}{k} = \frac{3.23 \times 1.6022 \times 10^{-19} \text{ J}}{1.38 \times 10^{-23} \text{ J K}^{-1}}$   
(assume 1 mole)  $= 3.75 \times 10^4 \text{ K}$   
(see also Table 4.11)

(c)  $f(E) = \frac{1}{e^{\frac{E-E_F}{kT}} + 1}$ ;  $\frac{(E-E_F)\text{eV} \times 1.6022 \times 10^{-19}}{1.38 \times 10^{-23} \times 300} = 38.7(E-E_F)$   
where  $E_F, E$  in eV.

$= \frac{1}{e^{38.7(E-E_F)} + 1}$

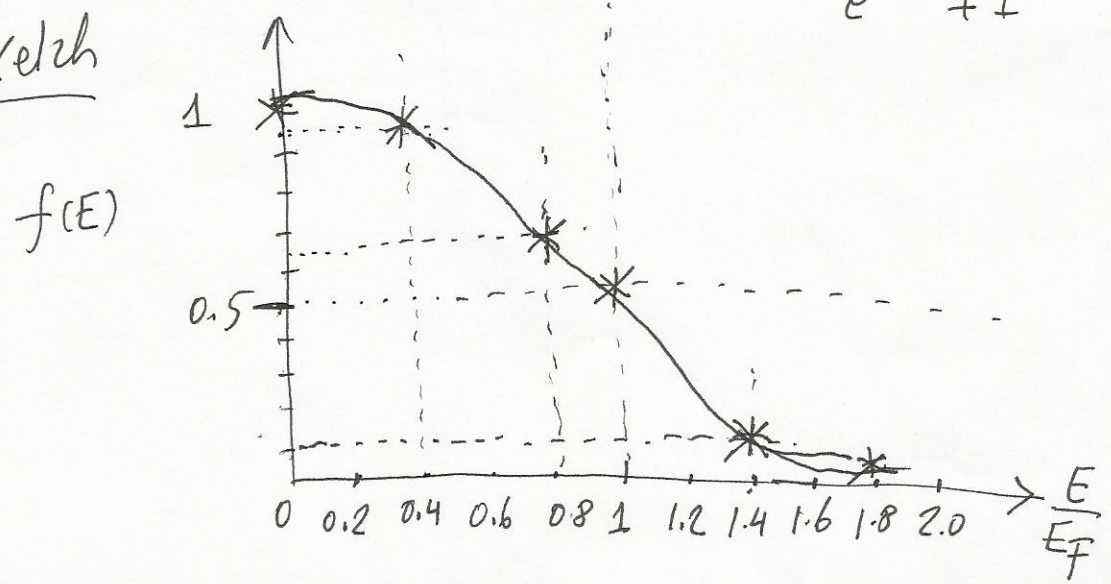
Table

$E/E_F$	0	0.4	0.8	1	1.4	1.8
$f(E)$	1	0.874	0.66	0.5	0.22	0.07

Example:  $\frac{E}{E_F} = 0.4 \Rightarrow E = 0.4 \times 3.23\text{eV} = 1.292\text{eV}$   
 $\therefore E - E_F = (1.292 - 3.23) = -1.938\text{eV}$

$\therefore f(0.4) = f(1.292) = \frac{1}{e^{\frac{-1.938}{kT}} + 1} = 0.874$

Sketch



$$(d) f(E) = \frac{1}{e^{38.7(E-E_F)} + 1}$$

↑  
eV

$$= \frac{1}{e^{38.7(E-3.23)} + 1}$$

$$\Rightarrow e^{38.7(E-3.23)} = f^{-1}(E) - 1$$

$$\Rightarrow 38.7(E-3.23) = \ln[f^{-1}(E) - 1]$$

$$\therefore E = \frac{\ln[f^{-1}(E) - 1]}{38.7} + 3.23$$

$f^{-1}(E)$	0.5	0.8	0.98
$E$	3.23 eV	3.19 eV	3.13 eV

no need to calculate

Example!  $f(E) = 0.8 \Rightarrow E = \frac{\ln\left[\frac{1}{0.8} - 1\right]}{38.7} + 3.23$

$$= \frac{\ln(0.25)}{38.7} + 3.23$$

$$= 3.19 \text{ eV}$$